

# Free Probability and Free Entropy

(A High-Level Introduction)

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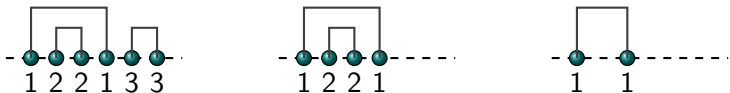
GOALS Summer School 2022

- 1 (Finishing Up) Proof of Free CLT
- 2 Microstates Free Entropy
- 3 Non-Microstates Free Entropy
- 4 Other Notions of Dimension
- 5 Takeaways

# What did you discover from Exercise 1?

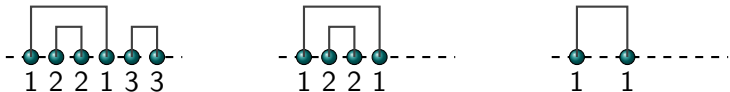
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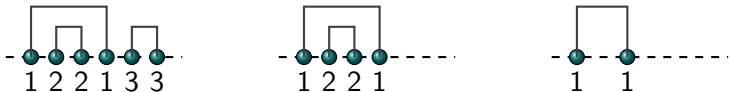
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Actually, this occurs if and only if  $\pi$  is non-crossing.

Otherwise,  $\phi(\pi) = 0$ .

# Finishing Up the Proof

At the end of the first lecture, we knew:

$$\lim_{k \rightarrow \infty} \phi(S_k^n) = \sum_{\pi \in \mathcal{P}_2(n)} \phi(\pi).$$

For even moments, since only non-crossing partitions  $\pi$  give a non-zero contribution, we have

$$\lim_{k \rightarrow \infty} \phi(S_k^{2n}) = \sum_{\pi \in NC_2(2n)} \phi(\pi) = \sigma^{2n} \cdot |NC_2(2n)|.$$

Since  $|NC_2(2n)| = C_n$ , the  $n$ th Catalan number, we are done!

## Theorem (Free Central Limit Theorem)

*If  $(a_i)_{i \in \mathbb{N}}$  are self-adjoint, freely independent, identically distributed  $nc$  random variables with  $\phi(a_i) = 0$  and  $\phi(a_i^2) = \sigma^2$ , then*

$$\frac{1}{\sqrt{k}}(a_1 + \cdots + a_k) = S_k \rightarrow \mathcal{S}(\sigma^2) \text{ in distribution.}$$

## Definition (Interpolated Free Group Factors)

Let  $(\mathcal{M}, \tau)$  be a tracial von Neumann algebra and let  $\mathcal{R}$  be a copy of the hyperfinite  $II_1$  factor in  $\mathcal{M}$ . Also let  $\omega = \{X^t \mid t \in T\}$  be a semicircular family such that  $\mathcal{R}$  and  $\omega$  are free.

Then for  $1 < r \leq \infty$ , we define  $L(\mathbb{F}_r)$  as the factor  $(\mathcal{R} \cup \{p_t X^t p_t \mid t \in T\})''$ , where  $p_t \in \mathcal{R}$  are projections satisfying  $r = 1 + \sum_{t \in T} \tau(p_t)^2$ .



## Aside: Compressions/Amplifications

For a  $II_1$  factor  $\mathcal{M}$  and  $0 < t < 1$ , we define the *compression* of  $\mathcal{M}$  as

$$\mathcal{M}_t \cong p\mathcal{M}p, \quad \text{for any } p \in \mathcal{P}(\mathcal{M}), \tau(p) = t.$$

We extend this notion to *amplifications* of  $\mathcal{M}$  by taking tensors with matrix algebras; for  $1 < t < \infty$ , if we write  $t = n \cdot \ell$ , where  $0 < \ell < 1$  and  $n \in \mathbb{N}$ , then we define

$$\mathcal{M}_t \cong p\mathcal{M}p \otimes M_n(\mathbb{C}) \cong M_n(p\mathcal{M}p), \quad \text{for any } p \in \mathcal{P}(\mathcal{M}), \tau(p) = \ell.$$

### Definition (Fundamental Group)

The **fundamental group** of a  $II_1$  factor  $\mathcal{M}$  is  $\{t \in \mathbb{R}_+ : \mathcal{M}_t \cong \mathcal{M}\}$ .  
It is a multiplicative subgroup of  $\mathbb{R}_+$ .

## Aside: Free Group Factor (FGF) Alternative

### Theorem (Dykema '92; Radulescu '94)

*One of the following two statements must be true:*

- 1  $L(\mathbb{F}_r) \cong L(\mathbb{F}_s)$  for all  $1 < r, s \leq \infty$ , and the fundamental group of  $L(\mathbb{F}_r)$  is  $\mathbb{R}_+$  for all  $1 < r \leq \infty$ .
- 2  $L(\mathbb{F}_r) \not\cong L(\mathbb{F}_s)$  for all  $1 < r, s \leq \infty$ , and the fundamental group of  $L(\mathbb{F}_r)$  is  $\{1\}$  for all  $1 < r \leq \infty$ .

### Addition Formula

$$L(\mathbb{F}_r) * L(\mathbb{F}_s) \cong L(\mathbb{F}_{r+s}), \text{ for } 1 < r, s \leq \infty.$$

### Compression Formula

$$L(\mathbb{F}_r)_t \cong L\left(\mathbb{F}\left(1 + \frac{r-1}{t^2}\right)\right), \text{ for } 1 < r \leq \infty, 0 < t < \infty.$$

# Microstates Free Entropy

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In classical information theory, entropy measures the amount of “information” or “uncertainty” in a random variable.

Boltzmann’s formula from physics says the entropy of a “macrostate” is obtained by counting how many “microstates” correspond to that “macrostate”.

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Classical entropy of a distribution  $\mu$  on  $\mathbb{R}^n$  with density  $\rho$  is given by

$$H(\mu) := - \int_{\mathbb{R}} \rho(x_1, \dots, x_n) \log \rho(x_1, \dots, x_n) dx_1 \cdots dx_n.$$

# Idea of $\chi(X)$

Let  $X = (x_1, \dots, x_n)$  be a tuple of self-adjoint elements of  $(M, \tau)$ .

Our “microstates” will be given by tuples of matrices in  $M_N(\mathbb{C})_{sa}$  that approximate the mixed moments of  $X$ .

“Counting the number of microstates” will be taking the Lebesgue measure in  $M_N(\mathbb{C})_{sa} \cong \mathbb{C}^{N^2} \cong \mathbb{R}^{2N^2}$ .

We then do some appropriate normalization and want to take limits as  $N \rightarrow \infty$ , and as the level of approximation improves.

# Definition of $\chi(X)$

Let  $(M, \tau)$  be a tracial  $W^*$ -probability space and let  $x_1, \dots, x_n$  be an  $n$ -tuple of self-adjoint elements in  $M$ . The set of approximating microstates is:

$$\begin{aligned} \Gamma(x_1, \dots, x_n; N, r, \epsilon) \\ := \{ (A_1, \dots, A_n) \in M_N(\mathbb{C})_{sa}^n : |\operatorname{tr}(A_{i_1} \cdots A_{i_k}) - \tau(x_{i_1} \cdots x_{i_k})| \leq \epsilon \\ \text{for all } 1 \leq i_1, \dots, i_k \leq n, 1 \leq k \leq r \}. \end{aligned}$$

Further define:

$$\chi(x_1, \dots, x_n; r, \epsilon) := \limsup_{N \rightarrow \infty} \left( \frac{1}{N^2} \log(\Lambda(\Gamma(x_1, \dots, x_n; N, r, \epsilon))) + \frac{n}{2} \log(N) \right).$$

The (microstates) free entropy  $\chi(x_1, \dots, x_n)$  is

$$\chi(x_1, \dots, x_n) := \lim_{\substack{r \rightarrow \infty \\ \epsilon \rightarrow 0}} \chi(x_1, \dots, x_n; r, \epsilon).$$



## Theorem

Let  $(M, \tau)$  be a tracial von Neumann algebra generated by self-adjoint  $x_1, \dots, x_n$ . Assume that  $\chi(x_1, \dots, x_n) > -\infty$ . Then

- 1 (Voiculescu '96)  $M$  does not have property  $\Gamma$ . In particular,  $M$  is a factor.

$M$  does not have property  $\Gamma$  if for any bounded sequence  $(t_k)$  s.t.

$\|[x, t_k]\|_2 \rightarrow 0$  for all  $x \in M$ , we have  $\|t_k - \tau(t_k)1\|_2 \rightarrow 0$  (every central sequence is trivial).

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- 3 (Ge '98)  $M$  is prime.

$M$  is prime if it cannot be decomposed as  $M = M_1 \bar{\otimes} M_2$  for  $II_1$  factors  $M_1, M_2$ .

# Non-Microstates Free Entropy

# Motivation

Another measure of the amount of “information” a random variable carries is the classical Fisher information.

Classically, entropy can be recovered through an appropriate integral of Fisher information.

This approach provides a more algebraic flavor and avoids the difficulty of finding the size of microstate spaces.

We will skip the classical formulation on this one!

# Non-Commutative Derivatives

## Definition

Define the *partial non-commutative derivatives*  $\partial_i$  as linear mappings

$$\partial_i : \mathbb{C} \langle X_1, \dots, X_n \rangle \rightarrow \mathbb{C} \langle X_1, \dots, X_n \rangle \otimes \mathbb{C} \langle X_1, \dots, X_n \rangle \text{ by}$$

$$\partial_i 1 = 0, \quad \partial_i X_j = \delta_{ij} 1 \otimes 1 \quad \text{for } j = 1, \dots, n,$$

and by the Leibniz rule:

$$\partial_i(P_1 P_2) = \partial_i(P_1) \cdot 1 \otimes P_2 + P_1 \otimes 1 \cdot \partial_i(P_2) \quad \text{for } P_1, P_2 \in \mathbb{C} \langle X_1, \dots, X_n \rangle.$$

**Exercise:** What is  $\partial_i$  on monomials? Compute

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**Exercise:** What is  $\partial_i$  on monomials? Compute

$$\partial_i(X_{i(1)} \cdots X_{i(m)}) = \sum_{k=1}^m \delta_{i,i(k)} X_{i(1)} \cdots X_{i(k-1)} \otimes X_{i(k+1)} \cdots X_{i(m)}.$$

## $\partial_i$ as operators on $L^2(M)$

Recall for  $x \in M$  we denote  $\|x\|_2^2 = \tau(x^*x)$ , and  $L^2(M)$  is the closure of  $M$  under this 2-norm.

If  $x_1, \dots, x_n \in M_{\text{sa}}$ , we can consider the operators  $\partial_i$  as derivatives on  $\mathbb{C}\langle x_1, \dots, x_n \rangle \subseteq M$  according to

$$\begin{array}{ccc} \mathbb{C}\langle X_1, \dots, X_n \rangle & \xrightarrow{\partial_i} & \mathbb{C}\langle X_1, \dots, X_n \rangle \otimes \mathbb{C}\langle X_1, \dots, X_n \rangle \\ \downarrow \text{eval} & & \downarrow \text{eval} \\ \mathbb{C}\langle x_1, \dots, x_n \rangle & \longrightarrow & \mathbb{C}\langle x_1, \dots, x_n \rangle \otimes \mathbb{C}\langle x_1, \dots, x_n \rangle \end{array}$$

(if the evaluation map is an algebra isomorphism).



# Motivation for $\xi_i$ , the conjugate variables

As an operator on  $L^2(x_1, \dots, x_n)$ ,  $\partial_i$  is unbounded, but we want them to be “nice”, i.e. closable.

$\implies \partial_i^*$  should be densely defined, i.e.  $1 \otimes 1 \in D(\partial_i^*)$ .

If  $1 \otimes 1 \in D(\partial_i^*)$ , then set  $\xi_i := \partial_i^*(1 \otimes 1)$ . Then we have the following relation:

$$\begin{aligned}\tau(\xi_i P(x_1, \dots, x_n)) &= \langle \partial_i^*(1 \otimes 1) P(\bar{x}), 1 \rangle \\ &= \langle 1 \otimes 1, \partial_i P(\bar{x}) \rangle \\ &= \tau \otimes \tau(\partial_i P(\bar{x})).\end{aligned}$$

# Free Fisher Information

Let  $x_1, \dots, x_n$  be self-adjoint elements of  $(M, \tau)$ .

- ① We say  $\xi_1, \dots, \xi_n$  form a **conjugate system** for  $x_1, \dots, x_n$  if for all  $i$ ,  $\xi_i \in L^2(x_1, \dots, x_n)$ , and they satisfy the **conjugate relations**: for all  $P \in \mathbb{C}\langle X_1, \dots, X_n \rangle$ ,

$$\tau(\xi_i P(x_1, \dots, x_n)) = \tau \otimes \tau((\partial_i P)(x_1, \dots, x_n)).$$

- ② The **free Fisher information** of  $x_1, \dots, x_n$  is defined by

$$\Phi^*(x_1, \dots, x_n) = \begin{cases} \sum_{i=1}^n \|\xi_i\|_2^2, & \text{if } \xi_1, \dots, \xi_n \text{ is a conjugate system} \\ & \text{for } x_1, \dots, x_n \\ +\infty, & \text{if no conjugate system exists.} \end{cases}$$

# Additivity of $\Phi^*$ is Equivalent to Freeness

## Theorem

Consider  $x_1, \dots, x_n, y_1, \dots, y_m$  self-adjoint elements of the same tracial von Neumann algebra  $(M, \tau)$ .

Then  $\{x_1, \dots, x_n\}$  and  $\{y_1, \dots, y_m\}$  are free if and only if

$$\Phi^*(x_1, \dots, x_n, y_1, \dots, y_m) = \Phi^*(x_1, \dots, x_n) + \Phi^*(y_1, \dots, y_m).$$

Proof uses a formulation in terms of cumulants characterizing when  $\{\xi_i\}$  is a conjugate system for  $\{x_i\}$ .

# Unification Problem

It was an open question for quite some time whether  $\chi(X) = \chi^*(X)$ . Since the result  $\text{MIP}^* = \text{RE}$  gave a negative answer to the Connes Embedding Problem (this means there exist  $\text{II}_1$  factors that do not embed into any ultrapower of  $\mathcal{R}$ ), we now know there are cases where  $\chi^*(X) < \chi(X)$ .

It is still open whether or not  $\chi(X) = \chi^*(X)$  on all  $\mathcal{R}^{\mathcal{U}}$ -embeddable tracial von Neumann algebras.

It is also open whether or not  $\chi(X)$  is a von Neumann algebra invariant, i.e. if  $W^*(X) = W^*(Y)$ , then is it true that  $\chi(X) = \chi(Y)$ ?

# Brief Mention: Other Notions of Dimension

- Hayes' 1-bounded entropy
  - defined via covering numbers of the microstate spaces.
  - known to be a von Neumann algebra invariant.
- Charlesworth and Nelson's free Stein dimension
  - algebraic flavor; defined by taking dimension of an appropriate space of derivations.
  - known to be a  $*$ -algebra invariant.
- Jekel's free entropy for types in model theory
  - generalizes Hayes' 1-bounded entropy to expressions that involve sup and inf in addition to the trace polynomials.
  - known to be a von Neumann algebra invariant.

# Takeaways

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- (Vaguely) What is free entropy, and why do people try to study it?
- Free Probability, Operator Algebras, Random Matrices



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- **There is more to explore!**

Thank you!

Hope you learned something from this! Ask me questions anytime!