

**Laplace Transform:
Convolutions and the Dirac Delta Function**

1. Use convolutions to find

$$\mathcal{L}^{-1} \left\{ \frac{1}{(s^2 + k^2)^2} \right\},$$

where $k \neq 0$ is a constant.

2. Find

$$\mathcal{L}^{-1} \left\{ \frac{G(s)}{(s-1)^2 + 1} \right\},$$

where $G(s)$ is the Laplace transform of a piecewise continuous function $g(t)$ of exponential order.

3. Find the solution to the initial value problem:

$$y'' + k^2 y = g(t); \quad y(0) = \alpha, \quad y'(0) = \beta,$$

where $k \neq 0$, α and β are constants, and $g(t)$ is a piecewise continuous function of exponential order.

Solve the following equations – all examples of **Volterra integral equations**:

4. $x(t) = 3 \cos t + 5 \int_0^t \sin(t - \tau)x(\tau) d\tau.$

7. $y(t) = 1 + \int_0^t (\tau - t)y(\tau) d\tau.$

5. $y(t) = t^3 + \int_0^t (t - \tau)y(\tau) d\tau.$

6. $y(t) = e^t + \int_0^t y(\tau) d\tau.$

8. $y(t) = \int_0^t (\tau - t)y(\tau) d\tau.$

9. Kirchhoff's second law in a series RLC circuit (consisting of a resistor, an inductor, and a capacitor) is:

$$L \frac{di}{dt} + Ri(t) + \frac{1}{C} \int_0^t i(\tau) d\tau = E(t),$$

where $E(t)$ is the voltage. Find the current $i(t)$ in such an RLC circuit, where

$$L = 0.1h; \quad R = 2\Omega; \quad C = 0.1f; \quad i(0) = 0,$$

and the voltage function is given by

$$E(t) = 120t(1 - u_1(t)).$$

10. $y' - 3y = \delta(t - 2); \quad y(0) = 0.$

11. $y'' + y = \delta(t - 2\pi); \quad y(0) = 0, \quad y'(0) = 1.$

12. $y'' + y = \delta\left(t - \frac{\pi}{2}\right) + \delta\left(t - \frac{3\pi}{2}\right); \quad y(0) = y'(0) = 0.$

13. $y'' + 2y' = \delta(t - 1); \quad y(0) = 0, \quad y'(0) = 1.$

14. $y'' + 4y' + 5y = \delta(t - 2\pi); \quad y(0) = y'(0) = 0.$

15. $y'' + 4y' + 13y = \delta(t - \pi) + \delta(t - 3\pi); \quad y(0) = 1, \quad y'(0) = 0.$

Laplace transforms of some basic functions

$$\begin{aligned} \mathcal{L}\{1\} &= \frac{1}{s}; \quad s > 0 & \mathcal{L}\{\sin(kt)\} &= \frac{k}{s^2 + k^2}; \quad s > 0 & \mathcal{L}\{\cosh(kt)\} &= \frac{s}{s^2 - k^2}; \quad s > |k| \\ \mathcal{L}\{t^n\} &= \frac{n!}{s^{n+1}}; \quad s > 0 & \mathcal{L}\{\cos(kt)\} &= \frac{s}{s^2 + k^2}; \quad s > 0 & \mathcal{L}\{u_a(t)\} &= \frac{e^{-as}}{s}; \quad s > 0 \\ \mathcal{L}\{e^{kt}\} &= \frac{1}{s - k}; \quad s > k & \mathcal{L}\{\sinh(kt)\} &= \frac{k}{s^2 - k^2}; \quad s > |k| & \mathcal{L}\{\delta(t - t_0)\} &= e^{-st_0} \\ & & & & \mathcal{L}\{\delta(t)\} &= 1 \end{aligned}$$

Inverse Laplace transforms of some basic functions

$$\begin{aligned} \mathcal{L}^{-1}\left\{\frac{1}{s}\right\} &= 1 & \mathcal{L}^{-1}\left\{\frac{1}{s^2 + k^2}\right\} &= \frac{1}{k} \sin(kt) & \mathcal{L}^{-1}\left\{\frac{s}{s^2 - k^2}\right\} &= \cosh(kt) \\ \mathcal{L}^{-1}\left\{\frac{1}{s^n}\right\} &= \frac{1}{(n-1)!} t^{n-1} & \mathcal{L}^{-1}\left\{\frac{s}{s^2 + k^2}\right\} &= \cos(kt) & \mathcal{L}^{-1}\left\{\frac{e^{-as}}{s}\right\} &= u_a(t) \\ \mathcal{L}^{-1}\left\{\frac{1}{s - k}\right\} &= e^{kt} & \mathcal{L}^{-1}\left\{\frac{1}{s^2 - k^2}\right\} &= \frac{1}{k} \sinh(kt) & \mathcal{L}^{-1}\{e^{-st_0}\} &= \delta(t - t_0) \\ & & & & \mathcal{L}^{-1}\{1\} &= \delta(t) \end{aligned}$$

Properties of the Laplace and Inverse Laplace transform

Translation Theorem I:

$$\begin{aligned} \mathcal{L}\{e^{kt} f(t)\} &= F(s - k) = \mathcal{L}\{f(t)\}|_{s \rightarrow s-k} \\ \mathcal{L}^{-1}\{F(s - k)\} &= \mathcal{L}^{-1}\{F(s)|_{s \rightarrow s-k}\} = e^{kt} \mathcal{L}^{-1}\{F(s)\} = e^{kt} f(t) \end{aligned}$$

Translation Theorem II:

$$\begin{aligned} \mathcal{L}\{f(t - a)u_a(t)\} &= e^{-as} F(s) = e^{-as} \mathcal{L}\{f(t)\} \\ \mathcal{L}^{-1}\{e^{-as} F(s)\} &= f(t - a)u_a(t) = \mathcal{L}^{-1}\{F(s)\}|_{t \rightarrow t-a} u_a(t) \end{aligned}$$

Derivatives of Laplace Transforms:

$$\begin{aligned} \mathcal{L}\{t^n f(t)\} &= (-1)^n \frac{d^n}{ds^n} F(s) \\ \mathcal{L}^{-1}\{F^{(n)}(s)\} &= (-1)^n t^n f(t) \end{aligned}$$

Laplace Transform of Periodic Functions:

If $f(t)$ is piecewise continuous on $[0, \infty)$, of exponential order, and periodic with period T :

$$\mathcal{L}\{f(t)\} = \frac{1}{1 - e^{-sT}} \int_0^T e^{-st} f(t) dt$$

Laplace Transforms of Derivatives:

$$\begin{aligned} \mathcal{L}\{y'\} &= sY(s) - y(0) \\ \mathcal{L}\{y''\} &= s^2 Y(s) - sy(0) - y'(0) \\ &\vdots \\ \mathcal{L}\{y^{(n)}(t)\} &= s^n Y(s) - s^{n-1}y(0) - s^{n-2}y'(0) - \dots - sy^{(n-2)}(0) - y^{(n-1)}(0) \end{aligned}$$

Convolution Theorem:

$$\begin{aligned} \mathcal{L}\{f(t) \star g(t)\} &= F(s)G(s) = \mathcal{L}\{f(t)\}\mathcal{L}\{g(t)\} \\ \mathcal{L}^{-1}\{F(s)G(s)\} &= f(t) \star g(t) \end{aligned}$$

Dirac Delta Function:

$$\begin{aligned} \int_0^\infty f(t)\delta(t - t_0) dt &= f(t_0) \\ (f \star \delta)(t) &= f(t) \end{aligned}$$

CONVOLUTION $f * g$

Example: $y'' + y = g(t)$; $y(0) = y'(0) = 0$.

Second-order linear equation with constant coefficients, with some input function $g(t)$, piecewise continuous and of exponential order on $t \in [0, \infty)$.

Solve using Variation of Parameters:

$$\textcircled{y_c} \quad m^2 + 1 = 0 \Rightarrow m = \pm i \Rightarrow y_c = C_1 \cos t + C_2 \sin t.$$

$$\textcircled{y_p} \quad W = \begin{vmatrix} \cos t & \sin t \\ -\sin t & \cos t \end{vmatrix} = 1.$$

$$W_1 = \begin{vmatrix} 0 & \sin t \\ g(t) & \cos t \end{vmatrix} = -\sin t g(t) \Rightarrow u_1' = -\sin t g(t) \Rightarrow u_1 = -\int_0^t \sin \tau g(\tau) d\tau \quad (\text{FTC II})$$

$$W_2 = \begin{vmatrix} \cos t & 0 \\ -\sin t & g(t) \end{vmatrix} = g(t) \cos t \Rightarrow u_2' = \cos t g(t) \Rightarrow u_2 = \int_0^t \cos \tau g(\tau) d\tau$$

$$\Rightarrow y = C_1 \cos t + C_2 \sin t - \cos t \int_0^t \sin \tau g(\tau) d\tau + \sin t \int_0^t \cos \tau g(\tau) d\tau$$

$$\Rightarrow y' = -C_1 \sin t + C_2 \cos t + \sin t \int_0^t \sin \tau g(\tau) d\tau - \cos t \sin t g(t) + \cos t \int_0^t \cos \tau g(\tau) d\tau + \sin t \cos t g(t)$$

$$\left. \begin{array}{l} y(0) = 0 \Rightarrow C_1 = 0 \\ y'(0) = 0 \Rightarrow C_2 = 0 \end{array} \right\} \Rightarrow y(t) = -\cos t \int_0^t \sin \tau g(\tau) d\tau + \sin t \int_0^t \cos \tau g(\tau) d\tau \\ = \int_0^t (\sin t \cos \tau - \sin \tau \cos t) g(\tau) d\tau \\ = \int_0^t \sin(t - \tau) g(\tau) d\tau$$

$$\Rightarrow \text{Solution: } y(t) = \int_0^t \sin(t - \tau) g(\tau) d\tau \quad \left(\sin(a - b) = \sin a \cos b - \sin b \cos a \right)$$

↳ convolution integral

(Typical form of the output/response $y(t)$ to an input $g(t)$ in a linear eqn. w/ constant coefficients).

DEFINITION: For piecewise continuous functions f and g on $[0, \infty)$, the convolution of f and g is the function denoted $f * g$, defined by:

$$(f * g)(t) = \int_0^t f(t-\tau)g(\tau) d\tau$$

$$= \int_0^t f(\tau)g(t-\tau) d\tau = (g * f)(t)$$

(Convolution is commutative).

Example: $f(t) = e^t$; $g(t) = \sin t$

$$(f * g)(t) = \int_0^t e^{t-\tau} \sin \tau d\tau = e^t \int_0^t e^{-\tau} \sin \tau d\tau$$

$$= e^t \frac{1}{2} (-e^{-\tau} \sin \tau - e^{-\tau} \cos \tau + 1)$$

$$(f * g)(t) = \frac{1}{2} (e^t - \sin t - \cos t)$$

$$I = \int_0^t e^{-\tau} \sin \tau d\tau = - \int_0^t (e^{-\tau})' \sin \tau d\tau$$

$$= -e^{-\tau} \sin \tau \Big|_0^t + \int_0^t e^{-\tau} \cos \tau d\tau$$

$$= -e^{-t} \sin t - \int_0^t (e^{-\tau})' \cos \tau d\tau$$

$$= -e^{-t} \sin t - e^{-\tau} \cos \tau \Big|_0^t - \underbrace{\int_0^t e^{-\tau} \sin \tau d\tau}_I$$

$$\Rightarrow 2I = -e^{-t} \sin t - e^{-t} \cos t + 1$$

Convolution Theorem: If $f(t), g(t)$ are piecewise continuous and of exponential order on $t \in [0, \infty)$, then:

$$\mathcal{L}\{f * g\} = F(s)G(s) = \mathcal{L}\{f(t)\} \mathcal{L}\{g(t)\}$$

(Useful because it allows one to find the Laplace transform of a convolution without having to compute the convolution function).

Inverse Form:

$$\mathcal{L}^{-1}\{F(s)G(s)\} = (f * g)(t)$$

(Useful because it allows one to find inverse Laplace transforms of products $F(s)G(s)$ which are otherwise difficult to handle).

Example: $\mathcal{L}\left\{\int_0^t e^{\tau} \sin(t-\tau) d\tau\right\}$

$$\begin{aligned} &= \mathcal{L}\{e^t * \sin t\} \\ &= \mathcal{L}\{e^t\} \mathcal{L}\{\sin t\} \\ &= \frac{1}{s-1} \frac{1}{s^2+1} = \frac{1}{(s-1)(s^2+1)} \end{aligned}$$

Example: $\mathcal{L}^{-1}\left\{\frac{1}{(s^2+4)^2}\right\} = \mathcal{L}^{-1}\{F(s)G(s)\}$, where $F(s) = G(s) = \frac{1}{s^2+4}$
 $= (f * g)(t) \quad \Rightarrow f(t) = g(t) = \frac{1}{2} \sin(2t)$

$$\Rightarrow (f * g)(t) = \int_0^t \frac{1}{2} \sin(2\tau) \frac{1}{2} \sin(2(t-\tau)) d\tau$$

$$= \frac{1}{4} \int_0^t \sin(2\tau) \sin(2t-2\tau) d\tau$$

$$\sin A \sin B = \frac{1}{2} (\cos(A-B) - \cos(A+B))$$

$$= \frac{1}{8} \int_0^t (\cos(4\tau-2t) - \cos(2t)) d\tau$$

$$= \frac{1}{8} \left(\frac{1}{4} \sin(4\tau-2t) \Big|_{\tau=0}^t - t \cos(2t) \right)$$

$$= \frac{1}{8} \left(\frac{1}{4} \sin(2t) - \frac{1}{4} \sin(-2t) - t \cos(2t) \right)$$

$$= \frac{1}{8} \left(\frac{1}{2} \sin(2t) - t \cos(2t) \right) = \frac{\sin(2t) - 2t \cos(2t)}{16}$$

Example: $4y'' + y = g(t)$; $y(0) = 3$; $y'(0) = -7$.

$$\mathcal{L}\{4y'' + y\} = \mathcal{L}\{g(t)\}$$

$$4(s^2 Y(s) - sy(0) - y'(0)) + Y(s) = G(s)$$

$$(4s^2 + 1)Y(s) - 12s + 28 = G(s) \Rightarrow Y(s) = \frac{12s - 28}{4s^2 + 1} + \frac{G(s)}{4s^2 + 1}$$
$$= \frac{3s - 7}{s^2 + 1/4} + \frac{1}{4} \frac{1}{s^2 + 1/4} G(s)$$

$$\Rightarrow Y(s) = \mathcal{L}^{-1}\left\{\frac{3s - 7}{s^2 + 1/4}\right\} + \frac{1}{4} \mathcal{L}^{-1}\left\{\frac{1}{s^2 + 1/4} G(s)\right\}$$

$$3\mathcal{L}^{-1}\left\{\frac{s}{s^2 + 1/4}\right\} - 7\mathcal{L}^{-1}\left\{\frac{1}{s^2 + 1/4}\right\}$$

$$= 3\cos(1/2t) - 14\sin(1/2t)$$

$$\mathcal{L}^{-1}\{F(s)G(s)\} = (f * g)(t)$$

$$\text{with } f(t) = 2\sin(1/2t)$$

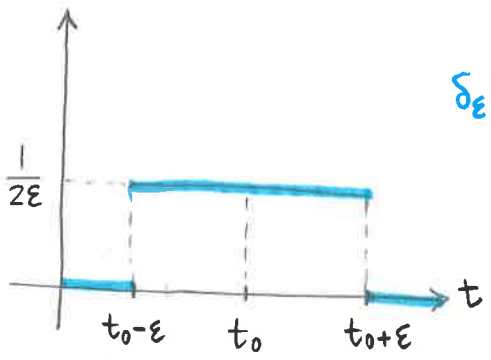
$$\Rightarrow Y(s) = 3\cos(1/2t) - 14\sin(1/2t) + \frac{1}{2}(\sin(1/2t) * g(t))$$
$$= 3\cos(1/2t) - 14\sin(1/2t) + \frac{1}{2} \int_0^t \sin\left(\frac{\tau}{2}\right) g(t-\tau) d\tau$$

Usefulness: This provides a general formula of the solution for an unspecified forcing function $g(t)$. So for a given $g(t)$, all that needs to be done is compute the convolution integral above.

The Dirac Delta Function

Unit Impulse: Model the situation of an external force of large magnitude acting on a system (mechanical, electrical, oscillatory etc.) only for a very short time. (a lightning strike; a golf ball given a sharp blow by a club etc.).

$$\delta_\epsilon(t-t_0) = \begin{cases} 0, & 0 \leq t < t_0 - \epsilon \\ \frac{1}{2\epsilon}, & t_0 - \epsilon \leq t < t_0 + \epsilon \\ 0, & t \geq t_0 + \epsilon \end{cases} = \frac{1}{2\epsilon} (u_{t_0-\epsilon}(t) - u_{t_0+\epsilon}(t))$$



$\delta_\epsilon(t-t_0)$: An impulse of magnitude $\frac{1}{2\epsilon}$, acting for (2ϵ) units of time, centered at $t=t_0$.

$$\int_0^\infty \delta_\epsilon(t-t_0) dt = \frac{1}{2\epsilon} (2\epsilon) = 1$$

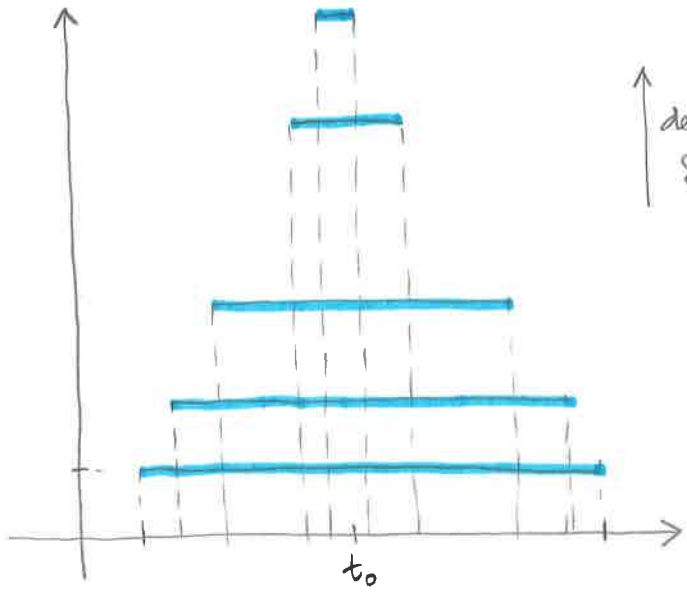
($\forall \epsilon > 0$): $\int_0^\infty \delta_\epsilon(t-t_0) dt = 1$ (therefore the name "unit" impulse).

Dirac Delta Function: To model a strong external force acting instantaneously at $t=t_0$

we define:

$$\delta(t-t_0) = \lim_{\epsilon \rightarrow 0} \delta_\epsilon(t-t_0) \quad \leftarrow \text{Dirac delta function}$$

Effectively, this lets the time interval $(t_0 - \epsilon, t_0 + \epsilon)$ around time t_0 decrease to 0. In the same time, the magnitude of the force, $\frac{1}{2\epsilon}$, tends to ∞ .



So in the limit, $\delta(t-t_0)$ is the "impulse".

$$\delta(t-t_0) = \begin{cases} 0, & \text{if } t \neq t_0 \\ \infty, & \text{if } t = t_0 \end{cases}$$

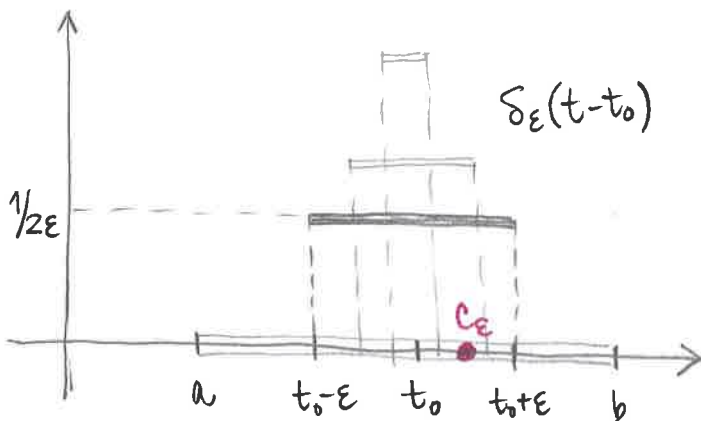
As you can see, this is not really a "function". In fact, rigorously speaking, δ is a distribution (aka generalized function), an advanced concept which is the subject of its own branch of mathematics - distribution theory. This began with Laurent Schwartz trying to make sense of Dirac's strange "function".

Roughly speaking, a distribution is an operator acting on functions, defined by how it acts on functions. For the Dirac delta, the defining action is:

For any continuous function f on an interval $[a, b]$ containing a point t_0 :

$$\int_a^b f(t) \delta(t-t_0) dt = f(t_0)$$

Proof:



$$\begin{aligned} \int_a^b f(t) \delta_\epsilon(t-t_0) dt &= \\ &= \int_{t_0-\epsilon}^{t_0+\epsilon} f(t) \frac{1}{2\epsilon} dt \\ &= \frac{1}{2\epsilon} (t_0+\epsilon - t_0-\epsilon) f(c_\epsilon) \\ &= f(c_\epsilon) \text{ for some } c_\epsilon \in [t_0-\epsilon, t_0+\epsilon] \end{aligned}$$

Mean Value Theorem for Integrals:

If f is continuous on $[a, b]$, then there is at least one value $c \in [a, b]$ such that:

$$\int_a^b f(x) dx = (b-a) f(c)$$

Letting $\epsilon \rightarrow 0$,

$$\int_a^b f(t) \delta(t-t_0) dt = f(t_0)$$

since $\lim_{\epsilon \rightarrow 0} f(c_\epsilon) = f(t_0)$
(this is forced by the interval $[t_0-\epsilon, t_0+\epsilon]$ collapsing onto the point $t=t_0$ as $\epsilon \rightarrow 0$).

Some Important Special Cases of:

$$\int_a^b f(t) \delta(t-t_0) dt = f(t_0)$$

$$a=0; b=\infty$$

$$\int_0^{\infty} f(t) \delta(t-t_0) dt = f(t_0)$$

aka the "sifting" property, because $\delta(t-t_0)$ "sifts" through all the values of f on $[0, \infty)$ to yield $f(t_0)$.

Taking $f(t) = e^{-st}$ above:

$$\int_0^{\infty} e^{-st} \delta(t-t_0) dt = e^{-st_0}$$

$\underbrace{\hspace{10em}}_{\mathcal{L}\{\delta(t-t_0)\}}$

$$\mathcal{L}\{\delta(t-t_0)\} = e^{-st_0}$$

In particular, letting $t_0=0$ above:

$$\mathcal{L}\{\delta(t)\} = 1$$

This is another indication that δ is not really a function: recall that for a piecewise continuous f , of exponential order, we expect $\lim_{s \rightarrow \infty} F(s) = 0$.

$$a=0; b=t_0$$

$$\int_0^{t_0} f(t) \delta(t_0-t) dt = f(t_0)$$

This follows because δ is even, i.e. $\delta(t-t_0) = \delta(t_0-t)$. But notice that this is really

$$(f * \delta)(t_0) = f(t_0).$$

In other words:

$$f * \delta = f$$

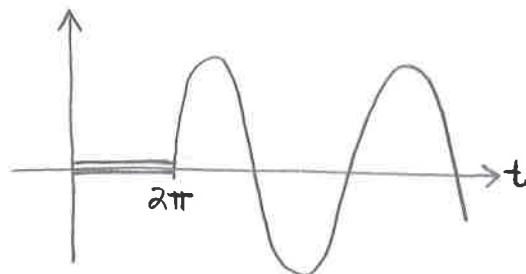
δ is the "identity" wrt convolution.

Example: $y'' + y = 4\delta(t - 2\pi)$ \rightarrow models the motion of a mass on a spring that is given a sharp blow at $t = 2\pi$.
 $y(0) = 0; y'(0) = 0$ \rightarrow mass begins at rest in the equilibrium position.

$$\mathcal{L}\{y'' + y\} = \mathcal{L}\{4\delta(t - 2\pi)\}$$

$$s^2 y(s) + y(s) = 4e^{-2\pi s} \Rightarrow y(s) = \frac{4e^{-2\pi s}}{s^2 + 1} \Rightarrow y(t) = 4\sin(t - 2\pi)u_{2\pi}(t) = 4\sin(t)u_{2\pi}(t)$$

$$y(t) = 4\sin(t)u_{2\pi}(t) = \begin{cases} 0, & \text{if } 0 \leq t < 2\pi \\ 4\sin t, & \text{if } t \geq 2\pi \end{cases}$$



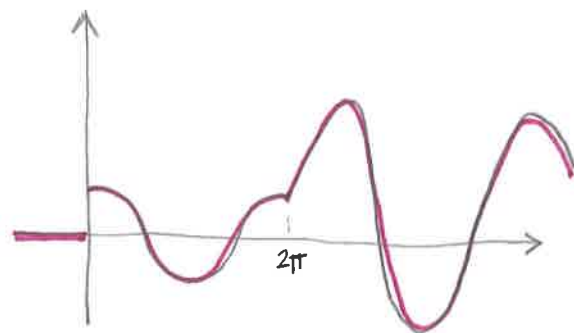
Example: $y'' + y = 4\delta(t - 2\pi)$
 $y(0) = 1; y'(0) = 0$ \rightarrow mass is released from rest 1 unit below the equilibrium point.

$$s^2 y(s) - s y(0) - y'(0) + y(s) = 4e^{-2\pi s}$$

$$(s^2 + 1)y(s) = s + 4e^{-2\pi s} \Rightarrow y(s) = \frac{s}{s^2 + 1} + \frac{4e^{-2\pi s}}{s^2 + 1}$$

$$\Rightarrow y(t) = \cos(t) + 4\sin(t)u_{2\pi}(t)$$

$$\Rightarrow y(t) = \begin{cases} \cos(t); & 0 \leq t < 2\pi \\ \cos(t) + 4\sin(t); & t \geq 2\pi \end{cases}$$



Remark: The condition $y'(0) = 0$ should be interpreted here in the sense that $\lim_{t \rightarrow 0^-} y'(t) = 0$.