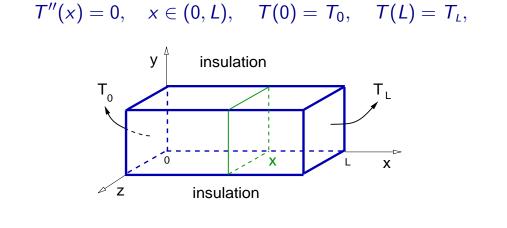


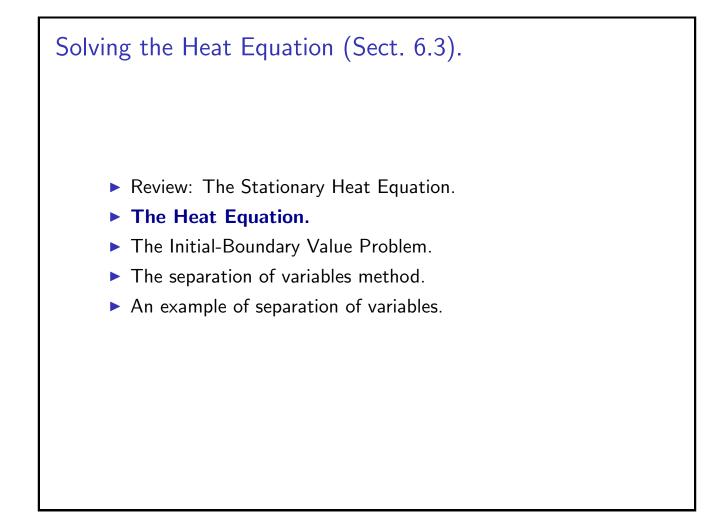
Review: The Stationary Heat Equation.

Review: The Stationary Heat Equation describes the temperature distribution in a solid material in thermal equilibrium. The temperature is time-independent.

Problem: The time-independent temperature, T, of a bar of length L with insulated horizontal sides and vertical extremes kept at fixed temperatures T_0 , T_L , is the solution of the BVP:



Remark: The heat transfer occurs only along the *x*-axis.



The Heat Equation.

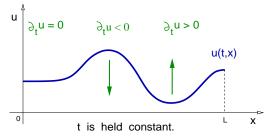
Remarks:

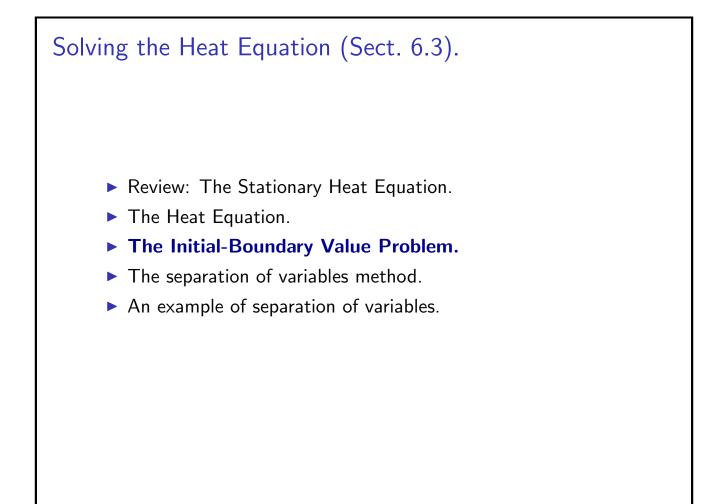
- The unknown of the problem is u(t, x), the temperature of the bar at the time t and position x.
- ▶ The temperature does not depend on *y* or *z*.
- ▶ The one-dimensional Heat Equation is:

$$\partial_t u(t,x) = k \, \partial_x^2 u(t,x),$$

where k > 0 is the heat conductivity, units: $[k] = \frac{(\text{distance})^2}{(\text{time})}$.

► The Heat Equation is a Partial Differential Equation, PDE.





The Initial-Boundary Value Problem.

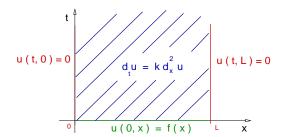
Definition

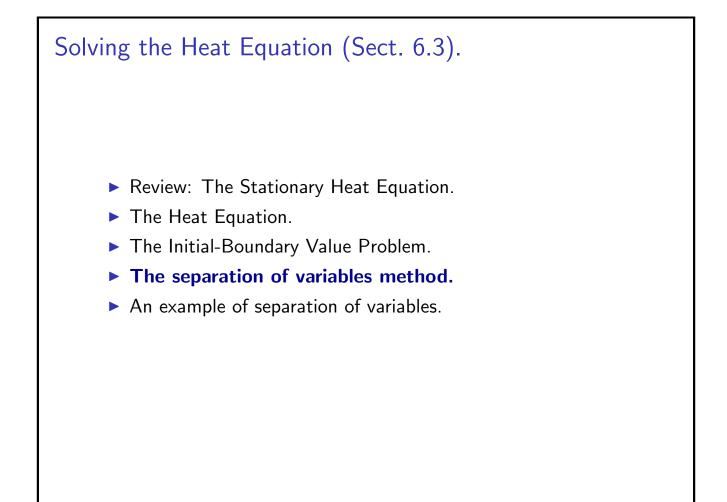
The IBVP for the one-dimensional Heat Equation is the following: Given a constant k > 0 and a function $f : [0, L] \to \mathbb{R}$ with f(0) = f(L) = 0, find $u : [0, \infty) \times [0, L] \to \mathbb{R}$ solution of

$$\partial_t u(t,x) = k \,\partial_x^2 u(t,x),$$

I.C.:
$$u(0, x) = f(x)$$
,

B.C.:
$$u(t, 0) = 0$$
, $u(t, L) = 0$.





Summary: IBVP for the Heat Equation.

Propose:

$$u(t,x)=\sum_{n=1}^{\infty}c_n\,v_n(t)\,w_n(x).$$

where

- \triangleright v_n : Solution of an IVP.
- w_n : Solution of a BVP, an eigenvalue-eigenfunction problem.
- ► *c*_n: Fourier Series coefficients.

Remark:

The separation of variables method does not work for every PDE.

Summary:

- ▶ The idea is to transform the PDE into infinitely many ODEs.
- We describe this method in 6 steps.

Step 1:

One looks for solutions u given by an infinite series of simpler functions, u_n , that is,

$$u(t,x)=\sum_{n=1}^{\infty}c_n\,u_n(t,x),$$

where u_n is simpler than u is the sense,

$$u_n(t,x) = v_n(t) w_n(x).$$

Here c_n are constants, $n = 1, 2, \cdots$.

The separation of variables method.

Step 2:

Introduce the series expansion for u into the Heat Equation,

$$\partial_t u - k \,\partial_x^2 u = 0 \quad \Rightarrow \quad \sum_{n=1}^{\infty} c_n \left[\partial_t u_n - k \,\partial_x^2 u_n \right] = 0.$$

A sufficient condition for the equation above is: To find u_n , for $n = 1, 2, \cdots$, solutions of

$$\partial_t u_n - k \, \partial_x^2 u_n = 0.$$

Step 3: Find $u_n(t,x) = v_n(t) w_n(x)$ solution of the IBVP

$$\partial_t u_n - k \partial_x^2 u_n = 0.$$

I.C.: $u_n(0, x) = w_n(x),$
B.C.: $u_n(t, 0) = 0, \quad u_n(t, L) = 0$

Step 4: (Key step.) Transform the IBVP for u_n into: (a) IVP for v_n ; (b) BVP for w_n . Notice: $\partial_t u_n(t,x) = \partial_t [v_n(t) w_n(x)] = w_n(x) \frac{dv_n}{dt}(t).$ $\partial_x^2 u_n(t,x) = \partial_x^2 [v_n(t) w_n(x)] = v_n(t) \frac{d^2 w_n}{dx^2}(x).$ Therefore, the equation $\partial_t u_n = k \partial_x^2 u_n$ is given by $w_n(x) \frac{dv_n}{dt}(t) = k v_n(t) \frac{d^2 w_n}{dx^2}(x).$ $\frac{1}{k v_n(t)} \frac{dv_n}{dt}(t) = \frac{1}{w_n(x)} \frac{d^2 w_n}{dx^2}(x).$ Depends only on t = Depends only on x.

The separation of variables method.

Recall:

$$\frac{1}{k v_n(t)} \frac{dv_n}{dt}(t) = \frac{1}{w_n(x)} \frac{d^2 w_n}{dx^2}(x).$$

Depends only on t = Depends only on x.

- The Heat Equation has the following property: The left-hand side depends only on t, while the right-hand side depends only on x.
- When this happens in a PDE, one can use the separation of variables method on that PDE.
- We conclude that for appropriate constants λ_m holds

$$\frac{1}{k v_n(t)} \frac{dv_n}{dt}(t) = -\lambda_n, \qquad \frac{1}{w_n(x)} \frac{d^2 w_n}{dx^2}(x) = -\lambda_n.$$

We have transformed the original PDE into infinitely many ODEs parametrized by n, positive integer. The separation of variables method. Summary Step 4: The original *IBVP* for the Heat Equation, PDE, can transformed into: (a) We choose to solve the following IVP for v_n , $\frac{1}{k v_n(t)} \frac{dv_n}{dt}(t) = -\lambda_n$, I.C.: $v_n(0) = 1$. Remark: This choice of I.C. simplifies the problem. (b) The BVP for w_n , $\frac{1}{w_n(x)} \frac{d^2w_n}{dx^2}(x) = -\lambda_n$, B.C.: $w_n(0) = 0$, $w_n(L) = 0$. Step 5: (a) Solve the IVP for v_n . (b) Solve the BVP for w_n .

The separation of variables method. Step 5(a): Solving the IVP for v_n . $v'_n(t) + k\lambda_n v_n(t) = 0$, I.C.: $v_n(0) = 1$. The integrating factor method implies that $\mu(t) = e^{k\lambda_n t}$. $e^{k\lambda_n t}v'_n(t) + k\lambda_n e^{k\lambda_n t} v_n(t) = 0 \Rightarrow [e^{k\lambda_n t}v_n(t)]' = 0$. $e^{k\lambda_n t}v_n(t) = c_n \Rightarrow v_n(t) = c_n e^{-k\lambda_n t}$. $1 = v_n(0) = c \Rightarrow v_n(t) = e^{-k\lambda_n t}$. The separation of variables method. Step 5(a): Recall: $v_n(t) = e^{-k\lambda_n t}$. Step 5(b): Eigenvalue-eigenvector problem for w_n : Find the eigenvalues λ_n and the non-zero eigenfunctions w_n solutions of the BVP $w_n''(x) + \lambda_n w_n(x) = 0$ B.C.: $w_n(0) = 0$, $w_n(L) = 0$. We know that this problem has solution only for $\lambda_n > 0$. Denote: $\lambda_n = \mu_n^2$. Proposing $w_n(x) = e^{r_n x}$, we get that $p(r_n) = r_n^2 + \mu_n^2 = 0 \implies r_{n\pm} = \pm \mu_n i$ The real-valued general solution is $w_n(x) = c_1 \cos(\mu_n x) + c_2 \sin(\mu_n x)$.

The separation of variables method.

Recall:
$$v_n(t) = e^{-k\lambda_n t}$$
, $w_n(x) = c_1 \cos(\mu_n x) + c_2 \sin(\mu_n x)$.

The boundary conditions imply,

$$0 = w_n(0) = c_1 \quad \Rightarrow \quad w_n(x) = c_2 \sin(\mu_n x).$$

$$0 = w_n(L) = c_2 \sin(\mu_n L), \quad c_2 \neq 0, \quad \Rightarrow \quad \sin(\mu_n L) = 0.$$
$$\mu_n L = n\pi \quad \Rightarrow \quad \mu_n = \frac{n\pi}{L} \quad \Rightarrow \quad \lambda_n = \left(\frac{n\pi}{L}\right)^2.$$

Choosing $c_2 = 1$, we get $w_n(x) = \sin\left(\frac{n\pi x}{L}\right)$.

We conclude that: $u_n(t,x) = e^{-k(\frac{n\pi}{L})^2 t} \sin(\frac{n\pi x}{L}), n = 1, 2, \cdots$

Step 6: Recall: $u_n(t,x) = e^{-k(\frac{n\pi}{L})^2 t} \sin(\frac{n\pi x}{L}).$

Compute the solution to the IBVP for the Heat Equation,

$$u(t,x) = \sum_{n=1}^{\infty} c_n u_n(t,x).$$
$$u(t,x) = \sum_{n=1}^{\infty} c_n e^{-k(\frac{n\pi}{L})^2 t} \sin\left(\frac{n\pi x}{L}\right).$$

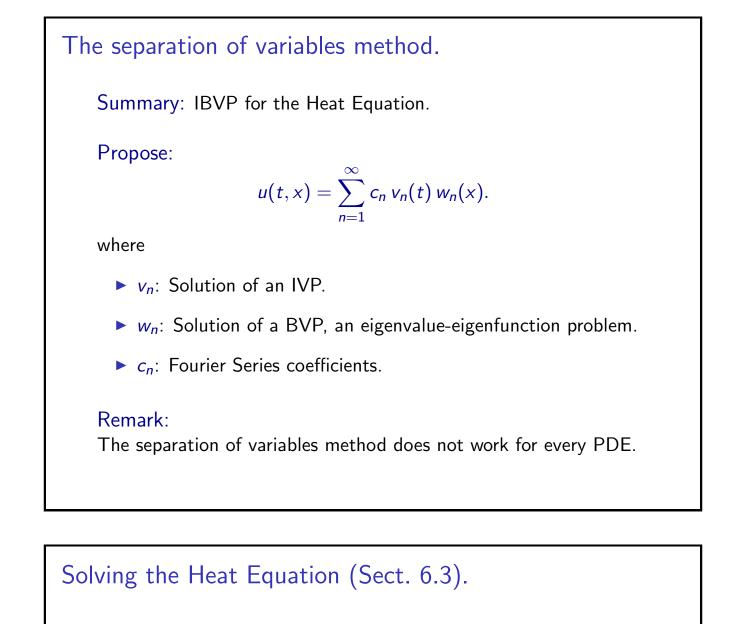
By construction, this solution satisfies the boundary conditions,

u(t,0) = 0, u(t,L) = 0.

Given a function f with f(0) = f(L) = 0, the solution u above satisfies the initial condition f(x) = u(0, x) iff holds

$$f(x) = \sum_{n=1}^{\infty} c_n \sin\left(\frac{n\pi x}{L}\right).$$

The separation of variables method. Recall: $u(t,x) = \sum_{n=1}^{\infty} c_n e^{-k(\frac{n\pi}{L})^2 t} \sin\left(\frac{n\pi x}{L}\right), \quad f(x) = \sum_{n=1}^{\infty} c_n \sin\left(\frac{n\pi x}{L}\right).$ This is a Sine Series for f. The coefficients c_n are computed in the usual way. Recall the orthogonality relation $\int_0^L \sin\left(\frac{n\pi x}{L}\right) \sin\left(\frac{m\pi x}{L}\right) dx = \begin{cases} 0, & m \neq n, \\ \frac{L}{2}, & m = n. \end{cases}$ Multiply the equation for u by $\sin\left(\frac{m\pi x}{L}\right)$ nd integrate, $\sum_{n=1}^{\infty} c_n \int_0^L \sin\left(\frac{n\pi x}{L}\right) \sin\left(\frac{m\pi x}{L}\right) dx = \int_0^L f(x) \sin\left(\frac{m\pi x}{L}\right) dx.$ $c_n = \frac{2}{L} \int_0^L f(x) \sin\left(\frac{n\pi x}{L}\right) dx, \quad u(t,x) = \sum_{n=1}^{\infty} c_n e^{-k(\frac{n\pi}{L})^2 t} \sin\left(\frac{n\pi x}{L}\right).$



- Review: The Stationary Heat Equation.
- ► The Heat Equation.
- ► The Initial-Boundary Value Problem.
- ► The separation of variables method.
- ► An example of separation of variables.

An example of separation of variables.

Example

Find the solution to the IBVP $4\partial_t u = \partial_x^2 u$, t > 0, $x \in [0, 2]$,

$$u(0,x) = 3\sin(\pi x/2), \quad u(t,0) = 0, \quad u(t,2) = 0.$$

Solution: Let $u_n(t,x) = v_n(t) w_n(x)$. Then

$$4w_n(x)\frac{dv}{dt}(t) = v_n(t)\frac{d^2w}{dx^2}(x) \quad \Rightarrow \quad \frac{4v'_n(t)}{v_n(t)} = \frac{w''_n(x)}{w_n(x)} = -\lambda_n.$$

The equations for v_n and w_n are

$$v'_n(t)+rac{\lambda_n}{4}v_n(t)=0,\qquad w''_n(x)+\lambda_n\,w_n(x)=0.$$

We solve for v_n with the initial condition $v_n(0) = 1$.

$$e^{\frac{\lambda_n}{4}t}v'_n(t)+\frac{\lambda_n}{4}e^{\frac{\lambda_n}{4}t}v_n(t)=0 \quad \Rightarrow \quad \left[e^{\frac{\lambda_n}{4}t}v_n(t)\right]'=0$$

An example of separation of variables.

Example

Find the solution to the IBVP $4\partial_t u = \partial_x^2 u$, t > 0, $x \in [0, 2]$,

$$u(0,x) = 3\sin(\pi x/2), \quad u(t,0) = 0, \quad u(t,2) = 0.$$

Solution: Recall: $\left[e^{\frac{\lambda_n}{4}t}v_n(t)\right]' = 0$. Therefore,

$$v_n(t)=c e^{-rac{\lambda_n}{4}t}, \quad 1=v_n(0)=c \quad \Rightarrow \quad v_n(t)=e^{-rac{\lambda_n}{4}t}.$$

Next the BVP: $w_n''(x) + \lambda_n w_n(x) = 0$, with $w_n(0) = w_n(L) = 0$. Since $\lambda_n > 0$, introduce $\lambda_n = \mu_n^2$. The characteristic polynomial is

$$p(r) = r^2 + \mu_n^2 = 0 \quad \Rightarrow \quad r_{n\pm} = \pm \mu_n i.$$

The general solution, $w_n(x) = c_1 \cos(\mu_n x) + c_2 \sin(\mu_n x)$.

The boundary conditions imply

$$0 = w_n(0) = c_1, \quad \Rightarrow \quad w_n(x) = c_2 \sin(\mu_n x).$$

An example of separation of variables.

Example

Find the solution to the IBVP $4\partial_t u = \partial_x^2 u$, t > 0, $x \in [0, 2]$, $u(0, x) = 3\sin(\pi x/2)$, u(t, 0) = 0, u(t, 2) = 0. Solution: Recall: $v_n(t) = e^{-\frac{\lambda_n}{4}t}$, and $w_n(x) = c_2 \sin(\mu_n x)$. $0 = w_n(2) = c_2 \sin(\mu_n 2)$, $c_2 \neq 0$, $\Rightarrow \sin(\mu_n 2) = 0$.

Then, $\mu_n 2 = n\pi$, that is, $\mu_n = \frac{n\pi}{2}$. Choosing $c_2 = 1$, we conclude,

$$\lambda_m = \left(\frac{n\pi}{2}\right)^2, \qquad w_n(x) = \sin\left(\frac{n\pi x}{2}\right)$$
$$u(t, x) = \sum_{n=1}^{\infty} c_n \, e^{-\left(\frac{n\pi}{4}\right)^2 t} \, \sin\left(\frac{n\pi x}{2}\right).$$

An example of separation of variables.

Example

Find the solution to the IBVP $4\partial_t u = \partial_x^2 u$, t > 0, $x \in [0, 2]$, $u(0, x) = 3\sin(\pi x/2)$, u(t, 0) = 0, u(t, 2) = 0. Solution: Recall: $u(t, x) = \sum_{n=1}^{\infty} c_n e^{-(\frac{n\pi}{4})^2 t} \sin\left(\frac{n\pi x}{2}\right)$. The initial condition is $3\sin\left(\frac{\pi x}{2}\right) = \sum_{n=1}^{\infty} c_n \sin\left(\frac{n\pi x}{2}\right)$. The orthogonality of the sine functions implies $3\int_0^2 \sin\left(\frac{\pi x}{2}\right) \sin\left(\frac{m\pi x}{2}\right) dx = \sum_{n=1}^{\infty} \int_0^2 \sin\left(\frac{n\pi x}{2}\right) \sin\left(\frac{m\pi x}{2}\right) dx$. If $m \neq 1$, then $0 = c_m \frac{2}{2}$, that is, $c_m = 0$ for $m \neq 1$. Therefore, $3\sin\left(\frac{\pi x}{2}\right) = c_1 \sin\left(\frac{\pi x}{2}\right) \Rightarrow c_1 = 3$.

An example of separation of variables.

Example

Find the solution to the IBVP $4\partial_t u = \partial_x^2 u$, t > 0, $x \in [0, 2]$, $u(0, x) = 3\sin(\pi x/2)$, u(t, 0) = 0, u(t, 2) = 0.

Solution: We conclude that

$$u(t,x) = 3 e^{-(\frac{\pi}{4})^2 t} \sin\left(\frac{\pi x}{2}\right).$$