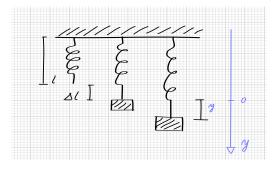


Consider a spring attached to the ceiling, having rest length I, with an attached mass m.

- (*l* + ∆*l*) is called equilibrium position of the spring loaded with a mass *m*.
- The coordinate y measures vertical deviations from the equilibrium position.



Forces acting on the system:

- Weight:  $F_g = mg$ .
- Spring:  $F_s = -k(\Delta l + y)$ . Hooke's Law. (Small oscillations.)
- Damping:  $F_d(t) = -d y'(t)$ . Fluid Resistance.

Newton's Law:  $my''(t) = F_g + F_s(t) + F_d(t)$ .

# Application: Mechanical Oscillations. Recall: $F_g = mg$ , $F_s = -k(\Delta l + y)$ , $F_d(t) = -dy'(t)$ . $my''(t) = F_g + F_s(t) + F_d(t)$ . That is, $my''(t) = mg - k(\Delta l + y(t)) - dy'(t)$ . At equilibrium, y = 0, y' = 0, then $k \Delta l = mg$ . Hence my''(t) = -ky(t) - dy'(t) my'' + dy' + ky = 0. To solve for the function y, we need the characteristic equation $mr^2 + dr + k = 0 \implies r_{\pm} = \frac{1}{2m} [-d \pm \sqrt{d^2 - 4mk}]$ .

Application: Mechanical Oscillations.

Recall: 
$$my'' + dy' + ky = 0$$
, and  $r_{\pm} = \frac{1}{2m} \left[ -d \pm \sqrt{d^2 - 4mk} \right]$ .

Not damped oscillations: d = 0. No fluid friction.

$$r_{\pm} = \pm \sqrt{-rac{k}{m}}, \qquad \omega_0 = \sqrt{rac{k}{m}}, \qquad r_{\pm} = \pm i\omega_0.$$

 $y(t) = c_1 \cos(\omega_0 t) + c_2 \sin(\omega_0 t).$ 

Remarks:

- Fundamental Frequency:  $\omega_0$ ; Period:  $T = \frac{2\pi}{\omega_0}$ .
- Equivalent expression:  $y(t) = A \cos(\omega_0 t \phi)$ .
- Amplitude: A; Phase shift:  $\phi$ .

Recall: Not damped oscillations:

$$y(t) = c_1 \cos(\omega_0 t) + c_2 \sin(\omega_0 t) \quad \Leftrightarrow \quad y(t) = A \cos(\omega_0 t - \phi).$$

where  $\omega_0 = \sqrt{k/m}$  is the fundamental frequency, A is the amplitude, and  $\phi$  the initial phase shift of the oscillations.

(Recall that the oscillation period is  $T = \frac{2\pi}{\omega_0}$ .)

**Proof:** Recall the trigonometric identity:

$$A\cos(\omega_0 t - \phi) = A\cos(\omega_0 t)\cos(\phi) + A\sin(\omega_0 t)\sin(\phi).$$

Therefore, comparing the first and last expressions above,

$$\begin{array}{l} c_1 = A\cos(\phi) \\ c_2 = A\sin(\phi) \end{array} \qquad \Leftrightarrow \qquad \begin{cases} A = \sqrt{c_1^2 + c_2^2} \\ \phi = \arctan\left(\frac{c_2}{c_1}\right). \end{array} \qquad \Box$$

Application: Mechanical Oscillations. Damped Oscillations Recall: my'' + dy' + ky = 0, and  $r_{\pm} = \frac{1}{2m} \left[ -d \pm \sqrt{d^2 - 4mk} \right]$ . Rewrite:  $r_{\pm} = -\frac{d}{2m} \pm \sqrt{\left(\frac{d}{2m}\right)^2 - \frac{k}{m}}$ . Introduce:  $\omega_0 = \sqrt{\frac{k}{m}}$ , and  $\omega_d = \frac{d}{2m}$ . Hence  $r_{\pm} = -\omega_d \pm \sqrt{\omega_d^2 - \omega_0^2}$ . Remark: We have three cases of damped oscillations:

(a) Over damped:  $\omega_d > \omega_0$ .

(b) Critically damped:  $\omega_d = \omega_0$ .

(c) Under damped:  $\omega_d < \omega_0$ .

Recall: m y'' + d y' + k y = 0, and  $r_{\pm} = -\omega_d \pm \sqrt{\omega_d^2 - \omega_0^2}$ .

(a) Over damped:  $\omega_d > \omega_0$ . Two distinct real roots:

$$y(t) = c_1 e^{r_+ t} + c_2 e^{r_- t}.$$

(b) Critically damped:  $\omega_d = \omega_0$ . Repeated real root  $r_{+} = r_{-} = \hat{r}$ :

$$y(t)=(c_1+c_2t)\,e^{\hat{r}t}.$$

(c) Under damped:  $\omega_d < \omega_0$ . Complex roots:

$$y(t) = [c_1 \cos(\beta t) + c_2 \sin(\beta t)] e^{-\omega_d t}$$

$$y(t) = A \cos(\beta t - \phi) e^{-\omega_d t}$$

where 
$$r_{\pm}=-\omega_{d}\pm ieta$$
, and  $eta=\sqrt{\omega_{0}^{2}-\omega_{d}^{2}}$ .

#### Application: Mechanical Oscillations.

#### Example

Find the movement of a 5Kg mass attached to a spring with constant  $k = 5\text{Kg/Secs}^2$  moving in a medium with damping constant d = 5Kg/Secs, with initial conditions  $y(0) = \sqrt{3}$  and y'(0) = 0.

Solution: The equation is: my'' + dy' + ky = 0, with m = 5, k = 5, d = 5. The characteristic roots are

$$r_{\pm}=-\omega_d\pm\sqrt{\omega_d^2-\omega_0^2},\quad \omega_d=rac{d}{2m}=rac{1}{2},\quad \omega_0=\sqrt{rac{k}{m}}=1.$$

 $r_{\pm} = -\frac{1}{2} \pm \sqrt{\frac{1}{4} - 1} = -\frac{1}{2} \pm i \frac{\sqrt{3}}{2}$ . Under damped oscillations.

$$y(t) = A \cos\left(\frac{\sqrt{3}}{2} t - \phi\right) e^{-t/2}.$$

#### Example

Find the movement of a 5Kg mass attached to a spring with constant  $k = 5\text{Kg/Secs}^2$  moving in a medium with damping constant d = 5Kg/Secs, with initial conditions  $y(0) = \sqrt{3}$  and y'(0) = 0.

Solution: Recall:  $y(t) = A \cos\left(\frac{\sqrt{3}}{2}t - \phi\right) e^{-t/2}$ . Hence,

$$y'(t) = -\frac{\sqrt{3}}{2}A\sin\left(\frac{\sqrt{3}}{2}t - \phi\right)e^{-t/2} - \frac{1}{2}A\cos\left(\frac{\sqrt{3}}{2}t - \phi\right)e^{-t/2}.$$

The initial conditions:

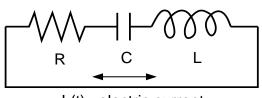
$$\sqrt{3} = y(0) = A\cos(\phi), \quad 0 = y'(0) = \frac{\sqrt{3}}{2}A\sin(\phi) - \frac{1}{2}A\cos(\phi).$$
$$\tan(\phi) = \frac{1}{\sqrt{3}} \quad \Rightarrow \quad \phi = \frac{\pi}{6}, \quad \Rightarrow \quad A = 2.$$
We conclude:  $y(t) = 2\cos\left(\frac{\sqrt{3}}{2}t - \frac{\pi}{6}\right)e^{-t/2}.$ 

Mechanical and electrical oscillations (Sect. 2.7?)

- Review: On solutions of  $y'' + a_1 y' + a_0 y = 0$ .
- Application: Mechanical Oscillations.
- ► Application: The RLC electrical circuit.

#### The RLC electrical circuit.

Consider an electric circuit with resistance R, non-zero capacitor C, and non-zero inductance L, as in the figure.



I (t) : electric current.

Kirchhoff's Law: The electric current flowing in the circuit satisfies:

$$L I'(t) + R I(t) + \frac{1}{C} \int_{t_0}^t I(s) ds = 0.$$

Derivate both sides above:  $LI''(t) + RI'(t) + \frac{1}{C}I(t) = 0.$ 

Divide by L:  $I''(t) + 2\left(\frac{R}{2L}\right)I'(t) + \frac{1}{LC}I(t) = 0.$ 

Introduce  $\alpha = \frac{R}{2L}$  and  $\omega = \frac{1}{\sqrt{LC}}$ , then  $I'' + 2\alpha I' + \omega^2 I = 0$ .

#### The RLC electrical circuit.

#### Example

Find real-valued fundamental solutions to  $I'' + 2\alpha I' + \omega^2 I = 0$ , where  $\alpha = R/(2L)$ ,  $\omega^2 = 1/(LC)$ , in the cases (a) (b) below.

Solution: The characteristic polynomial is  $p(r) = r^2 + 2\alpha r + \omega^2$ . The roots are:

$$r_{\pm} = \frac{1}{2} \left[ -2\alpha \pm \sqrt{4\alpha^2 - 4\omega^2} \right] \quad \Rightarrow \quad r_{\pm} = -\alpha \pm \sqrt{\alpha^2 - \omega^2}.$$

Case (a) R = 0. This implies  $\alpha = 0$ , so  $r_{\pm} = \pm i\omega$ . Therefore,

$$I_1(t) = \cos(\omega t), \qquad I_2(t) = \sin(\omega t).$$

Remark: When the circuit has no resistance, the current oscillates without dissipation.

