

Wednesday 12/04/13

Plan: * Review: Solving the Heat Equation

* Heat Eq.: Non-homogeneous B.C. (Not for Final Exam.)

* The Wave Equation. (Not for Final Exam.)

* Review: Solving the Heat Eq.

Example: Find a sol. of: $\begin{cases} \partial_t u = 2 \partial_x^2 u, & t \in [0, \infty), x \in [0, 1], \\ u(0, x) = f(x) = \begin{cases} 2 & x \in [0, \frac{1}{2}) \\ 0 & x \in [\frac{1}{2}, 1] \end{cases} \\ \partial_x u(t, 0) = 0, \quad u(t, 1) = 0. \end{cases}$

Sol: Decompose $u(t, x) = \sum_{n=1}^{\infty} c_n v_n(t) w_n(x)$ (2)

Introduce (2) into (1) $\sum_{n=1}^{\infty} \partial_t (c_n v_n w_n) = 2 \sum_{n=1}^{\infty} \partial_x^2 (c_n v_n w_n)$

$$\sum_{n=1}^{\infty} (\partial_t (c_n v_n w_n) - 2 \partial_x^2 (c_n v_n w_n)) = 0$$

$$\partial_t (c_n v_n w_n) = 2 \partial_x^2 (c_n v_n w_n)$$

$$c_n w_n(x) \dot{v}_n(t) = 2 c_n v_n(t) w_n''(x)$$

$$\boxed{\frac{1}{2} \frac{\dot{v}_n(t)}{v_n(t)} = \frac{w_n''(x)}{w_n(x)}}$$

$$\boxed{\begin{aligned} \frac{1}{2} \frac{\dot{v}_n(t)}{v_n(t)} &= -\lambda_n & \frac{w_n''(x)}{w_n(x)} &= -\lambda_n \\ v_n(0) &= 1 & w_n(0) &= 0 \\ v_n(t) &= e^{-2\lambda_n t} & w_n(1) &= 0 \\ \boxed{w_n(0) = 1} \end{aligned}}$$

$$v_n(t) = v_n(0) e^{-2\lambda_n t}$$

$$V_n(t) = e^{-2\lambda_n t}$$



$$V_n(t) = e^{-2((2n-1)\frac{\pi}{2})^2 t}$$

$$\begin{cases} W_n''(x) + \lambda_n W_n(x) = 0 \\ W_n'(0) = 0, \quad W_n(1) = 0 \end{cases}, \quad (W_n(0) = 1)$$

$$W_n(x) = c_1 \cos(\sqrt{\lambda_n} x) + c_2 \cancel{\sin(\sqrt{\lambda_n} x)}$$

$$W_n'(x) = -c_1 \sqrt{\lambda_n} \sin(\sqrt{\lambda_n} x) + c_2 \sqrt{\lambda_n} \cos(\sqrt{\lambda_n} x)$$

$$W_n'(0) = c_2 \sqrt{\lambda_n} = 0 \Rightarrow c_2 = 0$$

$$W_n(x) = c_1 \cos(\sqrt{\lambda_n} x)$$

$$W_n(1) = 0 = c_1 \cos(\sqrt{\lambda_n})$$

$$\downarrow \\ \sqrt{\lambda_n} = (2n-1) \frac{\pi}{2}$$

$$W_n(x) = c_1 \cos((2n-1) \frac{\pi}{2} x), \quad W_n(0) = 1 = c_1$$

$$W_n(x) = \cos((2n-1) \frac{\pi}{2} x)$$

$$u(t, x) = \sum_{n=1}^{\infty} c_n e^{-2((2n-1)\frac{\pi}{2})^2 t} \cos((2n-1) \frac{\pi}{2} x)$$

$$u(0, x) = f(x) = \sum_{n=1}^{\infty} c_n \cos((2n-1) \frac{\pi}{2} x)$$

$$\int_0^L f(x) \cos((2m-1) \frac{\pi}{2} x) dx = \sum_{n=1}^{\infty} c_n \int_0^L dx \cos((2n-1) \frac{\pi}{2} x) \cos((2m-1) \frac{\pi}{2} x)$$

$$\int_0^L f(x) \cos((2m-1) \frac{\pi}{2} x) dx = c_m \int_0^L \cos^2((2m-1) \frac{\pi}{2} x) dx = c_m \frac{L}{2}$$

$$c_m = \frac{2}{L} \int_0^L f(x) \cos((2m-1) \frac{\pi}{2} x) dx = \frac{2}{L} \int_0^{L/2} 2 \cos((2m-1) \frac{\pi}{2} x) dx$$

$$c_n = 4 \left(\frac{2}{(2n-1)\pi} \right) \sin((2n-1) \frac{\pi}{2} x) \Big|_0^{L/2}, \quad c_n = \frac{8}{(2n-1)\pi} \sin((2n-1) \frac{\pi}{4})$$

* Heat Eq. : Non-homogeneous B.C.

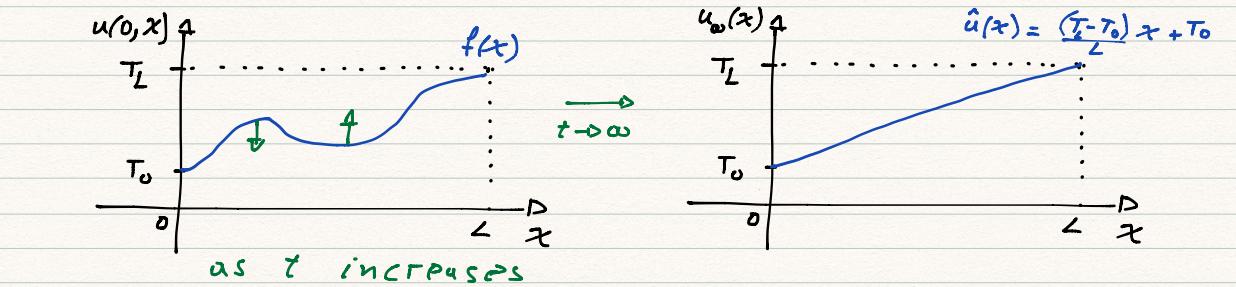
Consider the problem :

$$\left[\begin{array}{l} \text{Find } u \text{ sol. of: } \partial_t u = k \partial_x^2 u, \quad t \in [0, \infty), \quad x \in [0, L], \\ u(0, x) = f(x); \quad u(t, 0) = T_0, \quad u(t, L) = T_L \end{array} \right] \quad (1)$$

Now T_0, T_L might be non-zero.

Remark : A priori qualitative knowledge of the solution is needed to solve the problem.

- 1) Time evolution tries to smooth out high-curvature regions in the x -dependence of the solution.



Conclusion: There must exist a time-independent sol. to the Heat Eq. Let's call it u_∞ ; It is function of x .

u_∞ must be sol. of (1) in the limit $t \rightarrow \infty$

$$\partial_t u_\infty = k \partial_x^2 u_\infty, \quad \text{but } u_\infty \text{ depends only on } x.$$

$$u_\infty''(x) = 0; \quad \text{B.C. : } u_\infty(0) = T_0, \quad u_\infty(L) = T_L.$$

This problem has the unique sol:

$$\boxed{u_\infty(x) = \frac{(T_L - T_0)}{L} x + T_0}$$

- 2) Idea : Write : $u(t, x) = u_\infty(x) + \delta u(t, x)$

Find the problem for δu .

* Problem for δu :

$$\text{The Eq.: } \partial_t(u_\infty + \delta u) = k \partial_x^2(u_\infty + \delta u)$$

$$\begin{aligned} \partial_t u_\infty + \partial_t \delta u &= k \partial_x^2 u_\infty + k \partial_x \delta u \\ &= 0 & &= 0 \end{aligned}$$

$$\boxed{\partial_t(\delta u) = k \partial_x^2(\delta u)}$$

$$\text{The I.C.: } f(x) = u(0, x) = u_\infty(x) + \delta u(0, x)$$

$$u_\infty(x) = \frac{(T_L - T_0)}{L} x + T_0$$

$$\boxed{\delta u(0, x) = f(x) - \frac{(T_L - T_0)}{L} x - T_0}$$

$$\text{The B.C.: } T_0 = u(t, 0) = u_\infty(0) + \delta u(t, 0) = T_0 + \delta u(t, 0)$$

$$\boxed{\delta u(t, 0) = 0}$$

$$T_L = u(t, L) = u_\infty(L) + \delta u(t, L) = T_L + \delta u(t, L)$$

$$\boxed{\delta u(t, L) = 0}$$

Thrm: If u is sol. of $\partial_t u = k \partial_x^2 u$,
with I.C. $u(0, x) = f(x)$
B.C. $u(t, 0) = T_0$, $u(t, L) = T_L$

$$\text{Then } u(t, x) = u_\infty(x) + \delta u(t, x)$$

$$\text{with } u_\infty(x) = \frac{(T_L - T_0)}{L} x + T_0$$

and

$$\text{with I.C. } \partial_t \delta u = k \partial_x^2 \delta u$$

$$\delta u(0, x) = f(x) - u_\infty(x)$$

and B.C.

$$\delta u(t, 0) = 0, \quad \delta u(t, L) = 0.$$

* Solving the Wave Eq.

The IVP :

$$\text{Find } u \text{ sol. of: } \partial_t^2 u = a^2 \partial_x^2 u, \quad t \in [0, \infty), \quad x \in [0, L] \quad (1)$$

$$\text{I.C.} \quad u(0, x) = f(x), \quad \partial_t u(0, x) = g(x)$$

$$\text{B.C.} \quad u(t, 0) = 0, \quad u(t, L) = 0.$$

Sol.: consider the vector space:

$$V = \left\{ \tilde{v}, \text{ differentiable on } [0, L] \text{ s.t. } \tilde{v}(0) = \tilde{v}(L) = 0 \right\}.$$

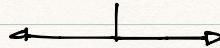
Let $\{w_n\}_{n=1}^{\infty}$ be an orthogonal basis of V

$$\text{Decompose: } u(t, x) = \sum_{n=1}^{\infty} v_n(t) w_n(x) \quad (2)$$

Introduce (2) in (1)

$$\sum_{n=1}^{\infty} \left[(\partial_t^2 v_n) w_n - a^2 v_n (\partial_x^2 w_n) \right] = 0 \Leftarrow (\partial_t^2 v_n) w_n = a^2 v_n (\partial_x^2 w_n)$$

$$\text{We focus on: } \frac{1}{a^2} \frac{\ddot{v}_n(t)}{v_n(t)} = \frac{w_n''(x)}{w_n(x)}$$



depends only on t

depends only on x

So each side must be a constant.

$$\boxed{\frac{1}{a^2} \frac{\ddot{v}_n(t)}{v_n(t)} = -\lambda_n} ; \quad \boxed{\frac{w_n''(x)}{w_n(x)} = -\lambda_n}$$

$w_n(0) = w_n(L) = 0.$

We start solving for w_n : $w_n(x) = \sin(\frac{n\pi}{2}x) ; \quad \lambda_n = \left(\frac{n\pi}{2}\right)^2$

$$\text{Recall: } \int_0^L \sin^2\left(\frac{n\pi}{2}x\right) dx = \frac{L}{2}$$

We now solve for V_n : $\dot{V}_n(t) + \alpha^2 \lambda_n V_n(t) = 0$

The general sol. is:
$$V_n(t) = c_n \cos(\alpha \frac{n\pi}{2} t) + d_n \sin(\alpha \frac{n\pi}{2} t)$$

Therefore: $u(t, x) = \sum_{n=1}^{\infty} [c_n \cos(\alpha \frac{n\pi}{2} t) + d_n \sin(\alpha \frac{n\pi}{2} t)] \sin(\frac{n\pi}{2} x)$

c_n, d_n are obtained from the I.C.s

$$f(x) = \sum_{n=1}^{\infty} c_n \sin\left(\frac{n\pi}{2} x\right)$$

$$g(x) = \sum_{n=1}^{\infty} d_n \frac{L}{\alpha n\pi} \sin\left(\frac{n\pi}{2} x\right)$$

Hence: $c_n = \frac{2}{L} \int_0^L f(x) \sin\left(\frac{n\pi}{2} x\right) dx$

$$d_n = \frac{2}{L} \alpha \frac{n\pi}{2} \int_0^L g(x) \sin\left(\frac{n\pi}{2} x\right) dx$$

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